

Grain Boundary Chloride Interface Segregation Enhances Conductivity in Hematite without Altering the Band Gap

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ABSTRACT

Nanocrystalline Fe₂O₃ with controlled chloride additions was used to evaluate how interfacial Cl segregation governs grain-boundary electrical transport. XRF quantified the retained chloride, while selective surface lixiviation and interface-excess analysis separated surface from grain-boundary (GB) contributions, revealing preferential segregation at GBs. X-ray diffraction showed hematite with modest crystallite-size changes. UV-Vis reflectance spectroscopy indicated an almost invariant optical band gap (~2.18 eV), whereas the color difference increased with retained Cl⁻, consistent with defect/surface states. Room-temperature impedance spectroscopy resolved bulk and GB responses and showed a marked drop in resistivity with increasing Cl⁻; after lixiviation, conductivity partially decreased, confirming the interfacial origin. The rise in

1 conductivity surpasses a simple carrier-count model ($\sigma \propto N$), implicating mobility
2 enhancement via GB-barrier lowering. Interfacial chloride segregation thus offers a
3 practical route to increase the electrical conductivity of polycrystalline Fe₂O₃ without
4 altering its band structure, advancing interface-engineered routes for electronic and
5 photoelectrochemical applications.
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12 Keywords: hematite (α -Fe₂O₃); chloride doping; interfacial segregation; grain-boundary
13 transport; impedance spectroscopy.
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21 **1. Introduction**

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24 Iron oxide (Fe₂O₃) nanoparticles are central to applications in catalysis, energy
25 conversion, and sensing [1-7], yet their performance is often limited by short carrier
26 lifetimes, low diffusion length, and intrinsically low electrical conductivity [8].
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29 Interface engineering, via interfacial segregation quantified by surface/grain boundary
30 interface excess and assisted by selective lixiviation, has emerged as an effective route
31 to mitigate these bottlenecks [9-11].
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38 A key pathway is interfacial segregation of dopants, which redistributes species between
39 surfaces and grain boundaries (GBs) and can lower the electrostatic barrier that blocks
40 charge transport across GBs [12-15]. This concept has improved stability and transport
41 in several oxides via quantified interface excess and targeted segregation, and recently
42 been demonstrated to tune grain-boundary conductivity in related ceramics [16].
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48 Notably, anion doping (e.g., halides) has been shown to enhance conductivity in
49 polycrystalline semiconductors by decreasing the GB barrier [10]. However, for
50 hematite, the role of chloride segregation, its partition between surface and GBs, its
51 persistence after selective lixiviation, and its direct impact on electrical transport, has
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1 not been systematically established or quantitatively connected to optical and
2 impedance responses.
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4 Here we test the hypothesis that retained chloride segregated at grain boundaries is the
5 primary driver for the conductivity enhancement in nanocrystalline Fe₂O₃, without
6 measurably altering the bulk band gap. We synthesize Fe₂O₃ with controlled chloride
7 additions, quantify the retained Cl⁻ by X-ray fluorescence (XRF), and use selective
8 lixiviation to separate surface from GB contributions [8]. The interface excess of Cl⁻ at
9 surfaces (Γ_s) and GBs (Γ_{gb}) is determined and related to its bandgap energy (E_g) and
10 room-temperature resistivity. Finally, we benchmark the conductivity increase against a
11 carrier-count model ($\sigma \propto N$) to demonstrate that the observed gains exceed what carrier
12 density alone can explain, implicating mobility enhancement via GB-barrier lowering.
13 This establishes interfacial Cl⁻ segregation as a practical lever to improve the electrical
14 performance of polycrystalline Fe₂O₃ while preserving its fundamental band structure.
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31 **2. Materials and methods**

32 **2.1. Synthesis of Cl-doped Fe₂O₃ nanopowders**

33 Fe₂O₃ nanopowders were synthesized by the polymeric precursor method [17]. The
34 synthesis used 45.8 wt % anhydrous citric acid (Synth, C₆H₈O₇, ≥ 99.00 wt % purity),
35 30.5 wt % ethylene glycol (Synth, C₂H₆O₂, ≥ 99.00 wt % purity), and 23.6 wt %
36 hydrated iron nitrate (Synth, Fe (NO₃)₃, ≥ 99.98 wt % purity) as the iron source.
37 Chloride was introduced by adding calculated amounts of NH₄Cl (Synth, ≥ 99.5 wt %
38 purity) aqueous solution to the resin. The target molar concentrations of Cl⁻ were x= 0.0,
39 0.1, 0.5, 1.0, 5.0, and 10.0 %. However, the molar concentrations of Cl⁻ obtained after
40 calcination were x= 0.0, 0.03, 0.04, 0.05, 0.18 and 0.23%. The lower amount of Cl⁻ was
41 probably due to evaporation losses during calcination. The Cl-doped samples were
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1 labeled as FeCl-x, and the undoped sample was labeled as Fe₂O₃. Calcination was
2 carried out under air flow at 300 °C for 10 h in a muffle furnace.
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7 **2.2. Powder characterization**

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9 Chemical compositions were determined by XRF using a Malvern Panalytical
10 spectrometer model Zetium. The X-ray diffraction (XRD) patterns were obtained using
11 an X'Pert-MPD (Philips) diffractometer with Cu K α radiation, operated at 45 kV and 40
12 mA. The step size was set at 0.02° 2 θ per second over the 5-90° 2 θ range. The
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2.3. Surface characterization

The selective lixiviation method was employed to determine the Cl⁻ surface excess [8].
Approximately 200 mg of powder was sonicated with ~1.8 g of distilled water for 1 h.
The matrix, *i.e.*, Fe₂O₃, which is insoluble in distilled water, was then separated from
the supernatant by centrifuging three times for 30 minutes at 13,000 rpm.
Approximately 1 g of the supernatant was collected, diluted in ~5 g of distilled water,
and analyzed using a chlorine selective electrode (Hanna Instruments).
Diffuse reflectance infrared Fourier transform (DRIFT) spectroscopy was performed
using a Nicolet iS50 spectrometer to identify the functional groups adsorbed on the

1 powder. DRIFT spectra were recorded in the range of 400-4000 cm^{-1} . Transmission
2 Electron Microscopy (TEM) images of Fe_2O_3 nanoparticles were obtained using a Titan
3 Themis Cubed (FEI Company) microscope equipped with double aberration correction
4 and operated at 300 kV.
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10 11 **2.4. Optical and electrical characterizations**

12 UV-vis spectra were investigated using a Shimadzu UV-2600 spectrometer equipped
13 with an integrating sphere to record the diffuse reflectance spectra (DRS) of the
14 samples.
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20 The color of the undoped Fe_2O_3 and Cl-doped Fe_2O_3 samples, both nonlixivated and
21 lixiviated, was evaluated using a MiniScan XE Plus colorimeter (HunterLab).
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24 Measurements were carried out according to the CIE Lab* color space standard (CIE
25 1976) [18], employing standard daylight (D65) as the illuminant and a 10° standard
26 observer. The instrument was calibrated using white and black reference tiles supplied
27 by the manufacturer.
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36 Impedance spectroscopy (IS) measurements were carried out using a Gamry reference
37 600+ potentiostat/galvanostat. Samples were prepared by dry pressing using a uniaxial
38 press at 294 MPa. Cylindrical pellets with 11 mm diameter and 1.5 mm thickness were
39 used in the experiments. Thermal treatment was performed at 325 $^\circ\text{C}$ for 2 h to enhance
40 the grain boundary area and enable electrical measurements. Low-temperature-curing
41 silver ink was used to deposit contact pads on the parallel surfaces of the pellets.
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51 Samples were dried at 80 $^\circ\text{C}$ under vacuum overnight before being assembled between
52 two stainless-steel current collectors in a 12 mm diameter Swagelok union-type cell
53 inside an Ar-filled glovebox (moisture and oxygen content below 2 ppm). IS data was
54 collected in the frequency range of 1 MHz to 0.01 Hz with 100 mV AC voltage
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1 amplitude at room temperature, within a Faraday cage (Gamry). The obtained spectra
2 were fitted using Z-View 4.0 software.
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7 **3. Results and Discussions**

8 **3.1. Powder characterizations**

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13 The X-ray diffraction patterns for undoped Fe₂O₃ and Cl-doped Fe₂O₃ samples are
14 shown in
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20 Figure 1. Broad peaks are observed, which is typical of nanosized crystallites. Hematite
21 (JCPDS card no. 24-0072) was identified as the main phase. A secondary maghemite
22 phase was detected only in the FeCl-0.18 sample, which is likely an effect of Cl⁻ doping.
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27 The average crystallite sizes and lattice parameters are shown in Table 1. The crystallite
28 size of the undoped Fe₂O₃ sample is consistent with the literature [19]. However, a
29 slight increase was observed in the doped samples up to 0.05 mol% Cl⁻, possibly due to
30 changes in diffusion parameters caused by Cl⁻ segregation at interfaces [20]. At low
31 doping levels, diffusion dominates, whereas at higher dopant contents interface
32 segregation induces the decrease in crystallite size. This decrease, observed for the
33 FeCl-0.18 and FeCl-0.23 samples, may be related to a balance in interfacial energy due
34 to segregation. Similar behavior has been reported for Sb-doped TiO₂ [21] and ZnO-
35 doped SnO₂ [22]. These two contributions to crystallite size evolution, diffusion and
36 surface energy, are more clearly illustrated by Eq. (1) [23, 24].
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$$54 \quad d_t^n - d_0^n = kD\gamma_{it} \quad (1)$$

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where, d_0 is the average particle size at t_0 ; D is the diffusion coefficient; γ_i is the interfacial energy; k is a constant that depends on the kinetic model; and n (typically 2, 3, or another integer) indicates the grain-growth mechanism [9].

Consistent with the behavior of FeCl-0.18 and FeCl-0.23, segregation lowers γ_i and, therefore, the growth rate predicted by Eq. (1). Consequently, reductions in interfacial energy, arising from changes in the interfacial chemical and from the build-up of interfacial excess (see Eqs. (1) and (2)), stabilize nanoparticles and inhibit grain-growth. This decrease in γ_i slows grain-growth kinetics [23, 24]. This nanoparticle stabilization is significant: the relative amount of segregation at grain boundaries versus surfaces governs the stability of these interfaces [24].

$$\gamma_i = \gamma_{i0} + \Gamma_i \Delta H_{seg} \quad \text{where } \Gamma_i = \frac{n}{A_i} \text{ and } \Delta H_{seg} < 0 \quad (2)$$

The lattice parameters of Fe₂O₃ were consistent with the literature [25, 26] and showed no significant changes upon Cl⁻ doping.

Sample	Hematite				Maghemite		
	Lattice parameters a=(b) (Å)	Lattice parameters c (Å)	Crystallite size (nm)	Phase fraction (wt%)	Lattice parameters a (Å)	Crystallite size (nm)	Phase fraction (wt%)
Fe ₂ O ₃	5.0377(5)	13.761(2)	37.57±0.78	100	-	-	-
FeCl-0.03	5.0377(5)	13.760(2)	39.52±0.73	100	-	-	-
FeCl-0.04	5.0376(4)	13.760(1)	46.57±1.07	100	-	-	-
FeCl-0.05	5.0372(4)	13.759(1)	46.88±1.22	100	-	-	-
FeCl-0.18	5.0368(4)	13.757(1)	41.20±0.85	86.6	8.353 (2)	27.14±2.49	13.4
FeCl-0.23	5.0381(4)	13.761(2)	40.22±0.88	100	-	-	-

Density, specific surface area (SSA), total specific area (TSA), specific grain-boundary (SGB) area, and SGB/SSA ratio are listed in Table 2. TSA was calculated using Eq. (3), assuming a truncated-octahedron crystallite shape, where d is the average crystallite size and ρ is the measured density [27]. The SGB area, defined as half the difference between TSA and SSA, was calculated using Eq. (4).

$$TSA = \frac{7.11}{d \cdot \rho} \quad (3)$$

$$SGB = \frac{(TSA - SSA)}{2} \quad (4)$$

The measured density decreased with increasing Cl⁻ content. This trend can be attributed to the formation of point defects, such as iron vacancies, generated to maintain electroneutrality when O²⁻ is replaced by Cl⁻ in the lattice [28].

Sample	Density (g/cm ³)	SSA (m ² /g)	TSA (m ² /g)	SGB (m ² /g)	SGB/SSA ratio
Fe ₂ O ₃	4.99±0.02	14.01±0.07	37.90±0.02	11.95±0.09	0.86±0.01
FeCl-0.03	4.99±0.01	13.89±0.03	35.99±0.02	11.10±0.08	0.80±0.01
FeCl-0.04	4.97±0.02	15.12±0.09	30.67±0.02	7.78±0.05	0.53±0.01
FeCl-0.05	4.97±0.02	14.09±0.09	30.49±0.03	8.20±0.06	0.56±0.01
FeCl-0.18	4.82±0.02	23.11±0.08	35.75±0.03	6.32±0.03	0.26±0.01
FeCl-0.23	4.83±0.03	27.74±0.16	36.55±0.02	4.40±0.02	0.14±0.01

At low additive levels (0.03-0.05 mol%), the SSA showed no significant change. By contrast, the FeCl-0.18 and FeCl-0.23 samples exhibited a slight increase in SSA. This

1 trend is consistent with a balance in interfacial energy driven by interfacial segregation
2 and is better interpreted together with SGB and the SGB/SSA ratio. For the Cl-doped
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4 Fe₂O₃ samples, both SGB and the SGB/SSA decreased, indicating that the dopant
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6 stabilizes surfaces more than grain boundaries, which also explains the higher SSA
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8 observed for FeCl-0.18 and FeCl-0.23. TSA decreased up to 0.05 mol% and then
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10 increased for FeCl-0.18 and FeCl-0.23 samples.
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17 **3.2. Interfacial segregation**

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19 For a better understanding of the effects of Cl⁻ doping in Fe₂O₃, it is important to
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21 quantify the dopant distribution at interfaces. The amount segregated at the nanoparticle
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23 surfaces was quantified directly by the selective lixiviation method, whereas dopant
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25 segregated at GBs or dissolved in the bulk was estimated. From the crystallite-size
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27 trend, the increase in particle size up to 0.05 mol% Cl⁻ is likely associated with the bulk
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29 solubility of Cl⁻ in Fe₂O₃ nanoparticles. Accordingly, in this study we considered the
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31 bulk solubility of Cl⁻ to be ~0.05 mol%. This estimate is consistent with the large
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33 differences between O²⁻ and Cl⁻ in ionic charge and radius, indicating low Cl⁻ solubility
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35 in the Fe₂O₃ matrix.
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41 Interfacial excess (Γ) was determined by selective lixiviation. First, the number of
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43 moles of Cl⁻ segregated at the grain boundaries (n_{gb}) was obtained as the difference
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45 between the total Cl⁻ content (n_T) and the amounts dissolved in the bulk (n_{bulk}) and
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47 segregated at the surface (n_S) as in Eq. (5) [23]:
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$$54 \quad n_{gb} = n_T - n_{bulk} - n_S \quad (5)$$

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59 Then, the grain-boundary excess (Γ_{gb}) was calculated using Eq. (6) [23]:
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$$\Gamma_{gb} = \frac{n_{gb}}{SGB} \quad (6)$$

The surface and grain-boundary excesses (Γ_S and Γ_{gb}) are summarized in **Error!**

Reference source not found..

Sample	Γ_S ($\mu\text{mol}/\text{m}^2$)	Γ_{gb} ($\mu\text{mol}/\text{m}^2$)
Fe ₂ O ₃	0.095±0.001	5.98±0.04
FeCl-0.03	0.179±0.001	5.98±0.06
FeCl-0.04	0.387±0.004	7.10±0.07
FeCl-0.05	0.651±0.005	9.40±0.08
FeCl-0.18	1.523±0.008	20.52±0.11
FeCl-0.23	1.973±0.013	21.45±0.14

Table 3 shows that most of the Cl⁻ segregated at the GBs rather than at particle surfaces.

The trace Cl⁻ detected in the undoped Fe₂O₃ sample most likely originated from precursor impurities in the Pechini method. This distribution is relevant because GB segregation is expected to lower the grain-boundary potential-barrier height for charge transport, thereby increasing the overall electrical conductivity of the system [10].

3.3. DRIFTS and TEM

DRIFTS and TEM were carried out on undoped and Cl-doped Fe₂O₃, in both lixiviated and nonlixiviated samples. The aim was to measure the effects of doping and of the selective lixiviation process.

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Figure 2a–c show the DRIFTS spectra of undoped and Cl-doped Fe₂O₃ before and after lixiviation. All spectra displayed well-defined bands assigned to adsorbed H₂O (3500–3200 cm⁻¹) [29], hydroxyl groups (~3680 and 1630 cm⁻¹) [29], adsorbed CO₂ (2330 cm⁻¹), adsorbed carbonate (CO₃²⁻, 1540–1290 cm⁻¹) and Fe–O bands (1000–400 cm⁻¹) [30, 31]. The features due to adsorbed H₂O, CO₂, and carbonate did not change upon lixiviation in any of the samples. Changes were observed only in the Fe–O region, particularly the Fe–OOH band, which can be attributed to the surface hydration during the selective lixiviation process.

The morphology and microstructure of the synthesized Fe₂O₃ samples were further characterized using TEM and high-resolution TEM (HRTEM). Figures 3a, 3c, and 3e present TEM micrographs of Fe₂O₃, FeCl-0.23, and lixivated FeCl-0.23, respectively.

The nanoparticles did not show any visible differences between the undoped and Cl-doped Fe₂O₃ samples. The same was observed for nonlixivated and lixivated samples (Figure 3c, e). These representative TEM micrographs show agglomerated nanoparticles forming elongated aggregates. Due to this irregular morphology and agglomeration, a reliable particle-size distribution could not be determined; nevertheless, spot measurements yielded particle sizes comparable to the crystallite size obtained from XRD. HRTEM lattice fringes confirmed well-crystallized Fe₂O₃ with a spacing of 0.37 nm, indexed to the (012) plane of hematite.

3.4. Optical and electrical characterizations

Figure 4 shows the UV-Vis diffuse reflectance spectra of undoped and Cl-doped Fe₂O₃. All samples exhibited a prominent reflectance peak in the 1.6–2.0 eV region, characteristic of hematite electronic transitions [32]. Increasing Cl⁻ content led to a progressive decrease in peak intensity, most pronounced for FeCl-0.18 and FeCl-0.23,

1 consistent with enhanced absorption in the visible range. Above ~2.2 eV the reflectance
 2 remained low for all compositions, with only minor variations among doping levels.
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 4 These trends are consistent with variations in defect/surface states rather than with a
 5
 6 modification of the bulk band structure, in agreement with the nearly constant band gap
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 8 obtained from Tauc plots (Figure 5). The optical band gap (E_g) was then determined
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 10 using the Tauc model [33], according to the Eq. (7):
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$$14 \quad (\alpha \cdot hv)^{\frac{1}{\beta}} = B(hv - E_g) \quad (7)$$

17 where α is the absorption coefficient dependent on photon energy, h is Planck's
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 19 constant, ν is the frequency of the incident radiation, and B is a material-dependent
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 21 constant. The value of β depends on the type of electronic transition, being 0.5 for direct
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 23 transitions and 2 for indirect transitions. α -Fe₂O₃ exhibits indirect electronic transitions
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 25 [34]; therefore, $\beta = 2$ was adopted in Eq. (7) to determine the band gap energy.
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28 To estimate α , the Kubelka–Munk function [35] was used, as given by Eq. (8), where
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 30 R_∞ represents the ratio between the reflectance of the sample and that of the standard:
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$$33 \quad \alpha \approx F(R) = \frac{(1-R_\infty)^2}{2R_\infty} \quad (8)$$

34 Substituting Eq. (8) into Eq. (7) yields the following relationship, which was used to
 35
 36 estimate the band gap energy:
 37
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$$39 \quad ((F(R) \cdot hv)^{\frac{1}{2}} = B(hv - E_g) \quad (9)$$

40 Figure 5 presents the corresponding Tauc plots for undoped and Cl-doped Fe₂O₃, from
 41
 42 which E_g values were extracted.
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45 The direct extrapolation of the linear region yields an optical band gap (E_g) of $\sim 2.16 \pm$
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 47 0.01 eV for all compositions. Thus, at the investigated concentrations, Cl⁻ does not
 48
 49 measurably change the fundamental band gap of hematite. The invariance of E_g
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1 indicates that chloride primarily influences the optical response through defect/surface
2 states, rather than by shifting the bulk band edges. Consistent with this interpretation,
3 the UV-vis reflectance shows a marked reduction in the 1.6–2.0 eV range (i.e.,
4 enhanced absorption), particularly for FeCl–0.18 and FeCl–0.23.
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10 Figure 6 presents the color difference ΔE^*_{ab} (CIE76) for undoped and Cl-doped Fe₂O₃,
11 referenced to the undoped sample under the same condition (nonlixivated or lixiviated),
12 as a function of the retained Cl⁻ content measured by XRF. ΔE was calculated using Eq.
13 (10). The individual L*, a*, and b* trends for each composition (before and after
14 lixiviation) are provided in the Supplementary Information (Figure 1S).
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$$26 \quad \Delta E = [(\Delta L^*)^2 + (\Delta a^*)^2 + (\Delta b^*)^2]^{1/2} \quad (10)$$

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30 ΔE^*_{ab} (CIE76) increases with the amount of retained Cl⁻ for both series, reaching visibly
31 distinct colors at higher contents; in the lixiviated series, a maximum is observed near
32 ~0.18 mol% and slightly declines at the highest composition. These trends indicate that
33 the perceived color is governed by the fraction of chloride effectively incorporated
34 and/or by defect states left after lixiviation, rather than by the nominal Cl⁻ added to the
35 sample. Consistently, the individual a* and b* coordinates show a net decrease with
36 increasing Cl (loss of red and yellow components), whereas L* varies less markedly;
37 the full L*, a*, b* evolutions before and after lixiviation are provided in the
38 Supplementary Information (Figure 1S).
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52 Guided by interfacial Cl⁻ segregation (Table 3), invariant E_g (Figure 5), and ΔE^*_{ab}
53 scaling with retained Cl⁻ (Figure 6), we hypothesize that chloride primarily affects
54 grain-boundary transport by lowering the interfacial barrier. To test this scenario, room-
55 temperature impedance spectra were acquired for undoped and Cl-doped Fe₂O₃ before
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1 (nonlixivated) and after lixiviation (Figure 7a, b). All samples exhibited similar
2 porosities ($\sim 50.0 \pm 8.0$ vol%) and were carefully prepared under identical experimental
3 conditions to minimize contributions from adsorbed species within the porous network.
4 This comparability allows differences in the IS response to be attributed primarily to
5 changes in the electrical microstructure (bulk vs. grain boundary), rather than to
6 morphological artifacts.
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8 The total electrical resistivity (ρ) of the undoped Fe_2O_3 pellet was $5.8 \times 10^8 \Omega\cdot\text{m}$ and
9 increased to $8.5 \times 10^8 \Omega\cdot\text{m}$ after lixiviation. Previous reports showed that $\alpha\text{-Fe}_2\text{O}_3$
10 exhibits a resistivity of $2 \times 10^6 \Omega\cdot\text{m}$ at 27°C [36], however, the authors did not report
11 the sample porosity.
12

13 All undoped and Cl-doped Fe_2O_3 samples exhibited two depressed semicircles and were
14 fitted with a series of R-CPE elements; in polycrystalline oxides, the higher-frequency
15 arc with lower capacitance is assigned to the bulk, whereas the lower-frequency arc with
16 higher capacitance corresponds to the grain boundary, in line with the brick-layer
17 description for electroceramics [37]. Any very-low-frequency contribution, when
18 present, is typically associated with electrode polarization. The total resistivity equals
19 the sum of the fitted resistive terms ($R_1 + R_2$), i.e., the difference between high- and
20 low-frequency Z' intercepts normalized by geometry.
21

22 With increasing Cl^- content, ρ decreases markedly for both series and partially recovers
23 after lixiviation, a trend consistent with grain-boundary barrier lowering driven by
24 interfacial chloride segregation. Similar behavior has been reported for Sn- and Zn-
25 doped hematite [12, 13], and analogous conductivity enhancements are observed in Cl-
26 doped ZnO nanoparticles processed under comparable conditions [8, 10]. As a
27 reference, undoped hematite is highly resistive at room temperature, and interface
28 engineering by dopants can shift its conductivity over several orders of magnitude [38].
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2 Figure 8 compiles the electrical conductivity (σ), calculated as $\sigma=1/\rho$, with ρ obtained
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4 from the total fitted resistance (R_1+R_2) and the sample geometry ($\rho=RA/L$) from the
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6 room-temperature IS data, as a function of the retained Cl^- content (XRF). Figure 8a
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8 compares nonlixivated and lixiviated pellets on common axes; Figure 8b shows the
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10 relative increase (σ/σ_0) together with the one-electron-per- Cl^- carrier-count model ($\sigma \propto$
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12 N), where σ_0 refers to the undoped sample under the same condition.
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16 From the combination of the lixiviation results and the conductivity trends in Figure 8a,
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18 two distinct contributions to the conductivity enhancement can be identified. The first is
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20 the ionic conduction provided by Cl^- species segregated at the surface of Fe_2O_3 , which
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22 strongly increases σ in the non-lixiviated pellets. Once these surface ions are removed
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24 by lixiviation, the conductivity drops sharply, revealing the second contribution: the
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26 electronic conduction facilitated by Cl^- remaining at grain boundaries, which lowers the
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28 grain-boundary barrier and thus sustains part of the conductivity. The reduced σ of the
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30 lixiviated samples therefore directly reflects the loss of the surface ionic pathway, while
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32 the residual conductivity is governed primarily by the grain-boundary electronic
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34 mechanism.
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38 Figure 8b benchmarks the conductivity gain against a simple carrier-count model
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40 derived from Eq. (10)

$$\sigma = N\mu|e| \quad (10)$$

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42 If each Cl^- contributed one electron and μ remained constant, the relative increase in σ
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44 should track the relative increase in N (taken as proportional to the retained Cl^-).
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48 Instead, the measured rise in σ far exceeds the change expected from N alone—most
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50 notably at higher Cl^- contents, indicating that the enhancement cannot be explained
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52 exclusively by carrier concentration. Put differently, to match the observed σ using Eq.
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1 (10) with a bulk-only picture would require effective carrier densities tens to hundreds
2 of times larger than those provided by the nominal Cl^- (as sketched by the “ N ” curve),
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4 which is unrealistic. We therefore infer a dominant increase in carrier mobility,
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6 consistent with grain-boundary barrier lowering driven by interfacial chloride; the
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8 partial reduction of σ after lixiviation further supports this interfacial mechanism [39].
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11 12 13 **4. Conclusions**

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16 This work shows that chloride introduced by the polymeric-precursor route modulates
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18 the interfacial chemistry of nanocrystalline Fe_2O_3 , driving preferential segregation at
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20 grain boundaries ($I_{\text{gb}} > I_{\text{s}}$). Selective lixiviation combined with interface-excess
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22 analysis separates surface and GB contributions, while XRD confirms hematite as the
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24 main phase with modest crystallite-size changes governed by the balance between
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26 diffusion and the reduction of interfacial energy induced by segregation.
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32 Optically, UV-vis DRS reveals an essentially invariant optical band gap (~ 2.18 eV); the
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34 color difference ΔE_{ab}^* increases with retained Cl^- , consistent with changes in
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36 defect/surface states rather than band-edge shifts. Room-temperature impedance
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38 spectroscopy resolves bulk and GB arcs and shows a marked drop in total resistivity
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40 with increasing Cl^- ; after lixiviation, the conductivity partially declines, confirming the
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42 interfacial origin of the effect. Moreover, the conductivity gain exceeds a simple carrier-
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44 count expectation ($\sigma \propto N$), indicating mobility enhancement via GB-barrier lowering
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46 rather than a bulk doping effect.
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52 Altogether, the data demonstrate that the conductivity improvement in Cl^- -modified
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54 Fe_2O_3 is predominantly governed by grain-boundary processes, achieved without
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56 altering the band structure. Interfacial Cl^- segregation therefore emerges as a practical
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1 route to tailor the electrical response of polycrystalline Fe₂O₃, with direct relevance to
2 electronic and photoelectrochemical applications.
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6

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7 Figures caption list
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11 Figure 1. X-ray diffraction patterns of undoped and Cl-doped Fe₂O₃ samples. Indexed
12 peaks correspond to hematite (JCPDS 24-0072); * indicates maghemite, observed only
13 in FeCl-0.18.
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19 Figure 2. DRIFTS spectra of undoped and Cl-doped Fe₂O₃ before and after lixiviation:
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21 (a) Fe₂O₃, (b) FeCl-0.18, (c) FeCl-0.23.
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24 Figure 3. TEM and high-resolution TEM (HRTEM) micrographs of Fe₂O₃, FeCl-0.23,
25 and lixiviated FeCl-0.23. (a,c,e) Low-magnification TEM (scale bar: 100 nm) showing
26 particle aggregates; yellow labels indicate apparent particle diameters. (b,d,f) HRTEM
27 (scale bar: 10 nm) with lattice fringes indexed to hematite, $d_{012} \approx 0.367$ nm (JCPDS 24-
28 0072). Panels correspond to: (a,b) Fe₂O₃, (c,d) FeCl-0.23, (e,f) FeCl-0.23 lixiviated.
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34 Figure 4. UV–Vis diffuse reflectance spectra of undoped and Cl-doped Fe₂O₃.
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38 Figure 5. Tauc plots used to determine the optical band gap (E_g) of undoped and Cl-
39 doped Fe₂O₃ samples: (a) Fe₂O₃, (b) FeCl-0.03, (c) FeCl-0.04, (d) FeCl-0.05, (e) FeCl-
40 0.18, and (f) FeCl-0.23.
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45 Figure 6. Color difference ΔE^*_{ab} (CIE76) for undoped and Cl-doped Fe₂O₃ as a
46 function of retained Cl⁻ content (XRF, mol%). ΔE was computed with respect to the
47 undoped sample under the same conditions (nonlixivated or lixiviated).
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52 Figure 7. Nyquist impedance plots ($-Z''$ vs Z') of undoped and Cl-doped Fe₂O₃ pellets
53 measured at room temperature: (a) nonlixivated, (b) lixiviated. Symbols are
54 experimental data and solid lines are the best fits to the equivalent circuit shown in the
55 inset (series of two R-CPE elements). Z' and $-Z''$ are expressed in $M\Omega \cdot m$.
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Figure 8. (a) Electrical conductivity, σ , versus retained Cl^- content (XRF, mol%) for nonlixivated and lixiviated pellets (room temperature). (b) Relative increase with respect to the undoped sample: measured σ (black squares) and carrier-count model N (red circles; one electron per Cl^-).

Figure 1S. Colorimetric parameters a^* (a), b^* (b), and L^* (c) of undoped Fe_2O_3 and Cl -doped Fe_2O_3 lixiviated and non-lixiviated samples.

Tables caption list

Table 1. Lattice parameters and crystallite sizes of undoped and Cl -doped Fe_2O_3 , and phase fractions obtained from Rietveld refinement. Values in parentheses are estimated standard deviations in the last digit. “ FeCl-x ” denotes the measured Cl content (mol%). Maghemite was detected only for FeCl-018 .

Table 2. Density, specific surface area (SSA), total specific surface area (TSA), specific grain-boundary area (SGB), and SGB/SSA ratio for undoped and Cl -doped Fe_2O_3 (FeCl-x). Values are reported as mean \pm standard uncertainty.

Table 3. Interfacial excess of Cl^- at the surface (Γ_s) and at the grain boundaries (Γ_{gb}) for undoped and Cl -doped Fe_2O_3 (FeCl-x) samples. Values are in $\mu\text{mol}/\text{m}^2$ and are reported as mean \pm standard uncertainty. excesses Γ_s is normalized by SSA; Γ_{gb} is normalized by SGB (see Eqs. (3)-(4) and (6)).















